

Travaux Dirigés de Physique Nucléaire et Physique des Particules

Shell model and applications

#1: Harmonic oscillator

1 / Explain why the energies of the 1D oscillator

$$h = \frac{P^2}{2\mu} + \frac{K}{2}X^2,$$

are proportional to $\hbar \sqrt{K/\mu}$. From now on, the strength is denoted $K = \mu \omega^2$.

2 / More explicitly, show that

$$h = \frac{\hbar \omega}{2} \tilde{h} = \frac{\hbar \omega}{2} \left[-\frac{d^2}{dx^2} + x^2 \right],$$

where \tilde{h} involves dimensionless quantities.

3 / Show that $\varphi_0(x) = \pi^{-1/4} \exp(-x^2/2)$ is the normalised ground-state of \tilde{h} with eigenvalue +1, either directly from the Schrödinger operator or from the lowering operator $a = (d/dx + x)/\sqrt{2}$.

4 / Deduce the two next eigenvalues and eigenstates of \tilde{h} using the raising operator $a^\dagger = (-d/dx + x)/\sqrt{2}$.

5 / (optional) Show that \tilde{h} has eigenvalues $1 + 2n$, where $n = 0, 1, 2, \dots$ and normalised eigenvectors

$$[2^n n! \sqrt{\pi}]^{-1/2} \exp(-x^2/2) H_n(x),$$

where H_n is an Hermite polynomial,¹ or for n even or odd, respectively

$$\left[\frac{n!}{\Gamma(n+3/2)} \right]^{1/2} \exp(-x^2/2) L_n^{-1/2}(x^2)$$

$$\left[\frac{n!}{\Gamma(n+3/2)} \right]^{1/2} \exp(-x^2/2) x L_n^{1/2}(x^2),$$

where L_n^β denotes an associated Laguerre polynomial, of which some useful properties are listed in Appendix.

6 / From the results on \tilde{h} , construct the two first s -wave levels and the first p -level of

$$H = \frac{\mathbf{p}^2}{2\mu} + \frac{\mu \omega^2}{2} \mathbf{r}^2.$$

¹The Hermite polynomials fulfill the orthogonality relation $\int_{-\infty}^{+\infty} \exp(-x^2) H_n(x) H_m(x) dx = \delta_{nm} 2^n n! \sqrt{\pi}$. $\exp(-x^2/2) H_n(x)$ obeys $y''(x) + (2n+1-x^2)y(x) = 0$.

Indicate the degeneracy of each level.

7 / H is now treated in spherical coordinates. Show that the eigenfunctions can be seek as

$$\Phi_{n\ell m}(\mathbf{r}) = R_{n\ell}(r) Y_\ell^m(\theta, \varphi)$$

where $Y_\ell^m(\theta, \varphi)$ is the usual spherical harmonics. Indicate the meaning of the quantum numbers n , ℓ and m . Explain why the radial function does not depend on m .

8 / (optional) Write the radial Schrödinger equation obeyed by $R_{n\ell}(r)$ and deduce

$$R_{n\ell}(r) = A (\alpha r)^\ell \exp[-\alpha^2 r^2/2] L_n^{\ell+1/2}(\alpha^2 r^2),$$

$$A = \left[\frac{2 n! \alpha^3}{\Gamma(n + \ell + 3/2)} \right]^{1/2},$$

9 / Show that the same energy levels $\hbar \omega(N+3/2)$ are obtained as by the method of Cartesian coordinates, with the same degeneracy $d(N)$. Are the eigenfunctions the same by the two methods?

#2: Proton density of ¹⁶O

1 / Show that for a spherical nucleus with closed shells, the proton density reads

$$\rho(r) = 2 \sum_{\substack{n,\ell \\ \text{occupied}}} \frac{2l+1}{4\pi} |R_{n\ell}(r)|^2$$

2 / Show that for the ground state of ¹⁶O, with 2 protons in the 1s shell and 6 in 1p, the charge density is

$$\rho(r) = \frac{2\alpha^3}{\pi^{3/2}} (1 + 2\alpha^2 r^2) \exp(-\alpha^2 r^2).$$

Plot the graph of this function $\rho(r)$.

#3: Energy levels, magic numbers

1 / Describe the energy spectrum corresponding to

- $V(r) = \mu\omega^2 r^2/2 + D \ell^2$
- $V(r) = \mu\omega^2 r^2/2 + D \ell^2 + g \ell \cdot s$

starting from a pure oscillator with energy levels $E = \hbar\omega(2n + \ell + 3/2)$, using $D = -0.025\hbar\omega$ and $g = -0.05\hbar\omega$. Describe the degeneracy for the pure oscillator and the above perturbed oscillators. Deduce the first magic numbers.

2 / Predict the spin and parity for the ground state, first excited state and one of the next excited states of the following nuclei: $^{209}_{82}\text{Pb}$, $^{209}_{83}\text{Bi}$, $^{41}_{20}\text{Ca}$. For both neutrons and protons, the shells are occupied in the following order:

$$\begin{aligned} &1s_{1/2} \quad 1p_{3/2} \quad 1p_{1/2} \quad 1d_{5/2} \quad 2s_{1/2} \quad 1d_{3/2} \\ &1f_{7/2} \quad 2p_{3/2} \quad 1f_{5/2} \quad 2p_{1/2} \quad 1g_{9/2} \\ &2d_{5/2} \quad 1g_{7/2} \quad 1h_{11/2} \quad 2d_{3/2} \quad 3s_{1/2} \\ &1h_{9/2} \quad 2f_{7/2} \quad 2f_{5/2} \quad 3p_{3/2} \quad 1i_{13/2} \quad 3p_{1/2} \\ &2g_{9/2} \quad 1i_{11/2} \quad 3d_{5/2} \quad 2g_{7/2} \dots \end{aligned}$$

#4: Schmidt lines

1 / Estimate the magnetic moment of a nucleus with an odd number A of nucleons, assuming it is due only to the last nucleon. The magnetic moment operator reads:

$$\boldsymbol{\mu} = \mu_N (g_s \mathbf{s} + g_\ell \boldsymbol{\ell}),$$

with $g_{sp} = 5.59$, $g_{sn} = -3.82$, $g_{lp} = 1$ et $g_{ln} = 0$ in units of the nuclear magneton, $e\hbar/(2m_p) = 3.15 \times 10^{-14}$ MeV/T, using the projection theorem

for a vector operator \mathbf{V} :

$$\langle \alpha jm | V_k | \alpha jm \rangle = \frac{\langle \alpha jm | \mathbf{V} \cdot \mathbf{J} | \alpha jm \rangle}{\langle \alpha jm | J^2 | \alpha jm \rangle} \times \langle \alpha jm | J_k | \alpha jm \rangle$$

2 / Estimate the expectation value of μ within the state $|jm\rangle$ with $m = j$.

3 / As an application, estimate the magnetic moment of the ground state of the nuclei $^{17}_9\text{F}$ and $^{17}_8\text{O}$.

Appendix: Laguerre polynomials

The associated Laguerre polynomials L_n^β are orthogonal polynomials, with increasing degree and increasing number of nodes, for the measure $\exp(-u)u^\beta$, in $u \in [0, \infty[$, with normalisation

$$\int_0^\infty \exp(-u) u^\beta L_n^\beta(u) L_m^\beta(u) du = \delta_{nm} \Gamma(\beta + n + 1)/n!.$$

Some examples are $L_0^{1/2}(u) = L_0^{3/2}(u) = 1$ and $L_1^{1/2}(u) = 3/2 - u$.

They obey, among others, the following differential equation

$$u y''(u) + (\beta + 1 - u) y'(u) + n y(u) = 0,$$

and it is straightforward to show that $\exp(-x^2/2) x^{\beta+1/2} L_n^\beta(x^2)$ obey

$$y''(x) + \left(4n + 2\beta + 2 - x^2 + \frac{1 - 4\beta^2}{4x^2} \right) y(x) = 0.$$

A generating function is

$$(1 - z)^{-\beta-1} \exp[xz/(z-1)] = \sum_n L_n^\beta(x) z^n.$$